## 三维样条小波单元构造及其在地下箱型结构抗爆 数值模拟中的应用

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摘要:由于材料的奇异性和加载的快速性,传统有限元在模拟爆炸荷载作用下地下结构动力响应过程中容易出现数值震荡,单元划分较多,计算效率低.小波有限元具有多尺度、多分辨率等特性,可以通过提高尺度函数阶数或小波函数的尺度来提高精度.用区间 B 样条尺度函数作为插值函数,推导了三维小波转换矩阵,构造了三维区间 B 样条小波单元,结合工程实例通过 Matlab 软件编程,对爆炸荷载作用地下箱形结构的动力响应进行了数值模拟,通过与 ANSYS/LS-DYNA 有限元程序模拟结果进行对比,小波有限元用较少的单元获得了较高的精度,提高了计算效率,避免了数值震荡.

**关键词:** 爆炸荷载; 地下箱形结构; 动力响应; 区间 B 样条小波, 单元构造

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用传统有限元对 6 m 跨地下支护结构进行爆炸荷载作用下的模拟计算时,计算结果出现了数值震荡,如图 1. 主要是因为爆炸荷载的加载速度非常快,加载梯度增大,荷载突增. 此外,混凝土与岩石的应力应变关系具有奇异性,这些给传统有限元带来了无法解决的问题.

小波有限元是将小波与传统有限元相结合,利用小波函数或尺度函数作为插值函数构造单元,使问题的求解在一个嵌套序列中进行,利用小波多分辨率的特性,在求解变化平缓区域采用较小的分析尺度以提高分析效率.由于小波函数具有多尺度、多分辨率与紧支性等特性,小波有限元是近年来得到迅速发展的一种数值分析方法,可以解决有限元分析中最难以处理的工程奇异性问题[1],样条函数是80年代发展起来

的现代函数逼近的一个重要的分支<sup>[2-3]</sup>,是现代数值计算的 重要基础,具有连续性强、待定系数少、高精度和高计算 效率等优点,目前为止已经具有完备的计算理论,所以样 条小波是将样条函数与小波理论相结合的产物,集样条函 数逼近精度、效率高和小波多分辩特性等优点于一体,开 辟了数值计算的新领域。由于一个区间 B 样条小波单元能 够求出很多个内部和边界节点值,同时,小波单元具有良 好的数值逼近性,因此在计算中可以用较少的单元获得较 高的精度。

国内外学者对样条小波做了大量研究. Chen W H 将小波与有限元相结合,构造了样条小波单元<sup>[1]</sup>,美国学者 Chui和 Quak构造了[0,1]区间上的 B 样条小波,生成了有限区间上的多分辨分析,并给出了快速分解和重构算法. 何正嘉、向家伟、陈雪峰等对区间 B 样条小波一维、二维单元的构造进行了详细研究<sup>[4-7]</sup>.

岩石地下结构地质条件复杂,爆炸荷载作用下地下结 构的动力响应,是军事地下工程重要的研究课题,也是一

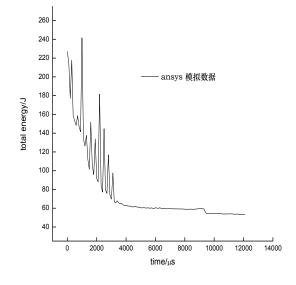


图 1 总的能量曲线 Fig. 1 The total energy curve

个非常复杂的力学变化过程,其中包含了多重非线性,目前,国内外应用小波有限元进行数值模拟研究的 文献很少,本文推导了小波三维转换矩阵,构造了三维实体小波单元,并通过 Matlab 编程将其应用在爆炸 荷载作用下的地下箱型结构的动力响应模拟中,很好的解决了数值震荡问题.

## 1 三维区间 B 样条小波单元构造

### 1.1 三维小波转换矩阵推导

三维区间样条六面体单元如图 2 所示,根据文献 1 的求解思路. 单元边值问题用  $L(u(x,y,z))=f(x,y,z)\in\Omega$  表示. 对六面体  $\Omega$  求解域进行离散化,将  $\Omega$  求解域映射为  $\Omega_e=\left\{\xi,\eta,\zeta\mid\xi,\eta,\zeta\in[0,1]\right\}$ . 采用三维张量积区间 B 样条函数尺度函数为插值函数,求解域划分为  $n\times n\times n$  个网格,节点数为 $(n+1)\times(n+1)\times(n+1)$  个. 则未知场函数 $u(\xi,\eta,\zeta)$  可以用式 (1)表示

$$u(\xi,\eta,\zeta) = \Phi a^e \tag{1}$$

式中, $\Phi$ 为小波插值基函数.

$$\boldsymbol{\Phi} = \boldsymbol{\Phi}_1 \otimes \boldsymbol{\Phi}_2 \otimes \boldsymbol{\Phi}_3 \tag{2}$$

式中:  $\Phi_1, \Phi_2,$ 和  $\Phi_3$  分别为 m 阶 j 尺度下的一维小波尺度函数.  $a^e$  为待求的小波插值系数列向量.

令  $\xi_i = (i-1)/n$ ;  $\eta_i = (i-1)/n$ ;  $\zeta_i = (i-1)/n$ ;  $i, j = 1, 2, \dots, n+1$  为标准求解域中各节点的坐标值,

设物理自由度列向量: 
$$u^e = [u_1^e \cdots u_{i1}^e \cdots u_{n+1}^e]^T$$
 (3)

可以得到: 
$$u^e = R^e a^e \tag{4}$$

式中: 
$$R^e = R_1^e \times R_2^e \times R_3^e \tag{5}$$

其中: 
$$\begin{cases} R_{1}^{e} = [\boldsymbol{\Phi}_{1}^{T}(\xi_{1})\boldsymbol{\Phi}_{1}^{T}(\xi_{2})\cdots\boldsymbol{\Phi}_{1}^{T}(\xi_{n+1})]^{T} \\ R_{2}^{e} = [\boldsymbol{\Phi}_{2}^{T}(\eta_{1})\boldsymbol{\Phi}_{2}^{T}(\eta_{2})\cdots\boldsymbol{\Phi}_{2}^{T}(\eta_{n+1})]^{T} \\ R_{3}^{e} = [\boldsymbol{\Phi}_{3}^{T}(\zeta)\boldsymbol{\Phi}_{3}^{T}(\zeta_{2})\cdots\boldsymbol{\Phi}_{3}^{T}(\zeta_{n+1})]^{T} \end{cases}$$
(6)

由公式 (1) 和公式 (4),可得到三维  $C_0$  型转换矩阵:

$$T^e = (R^e)^{-1} (8)$$

形函数  $N^e = \Phi T^e$  (9)

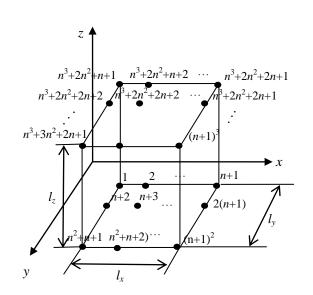


图 2 三维样条单元 Fig.2 Three-way interval B-spline sector wavelet element

#### 1.2 三维区间 B 样条小波单元构造

## 1.2.1 能量泛函分析

对于爆炸强动荷载作用下的结构分析,必须考虑惯性力的动力效应,设 c 为阻尼粘滞系数,  $\rho$  为拱的密度,三维几何方程为<sup>[8]</sup>

$$\varepsilon = LV = \begin{bmatrix} \frac{\partial}{\partial x} & & & \\ & \frac{\partial}{\partial y} & & \\ & & \frac{\partial}{\partial z} & \\ \frac{\partial}{\partial y} & \frac{\partial}{\partial x} & & \\ & \frac{\partial}{\partial z} & \frac{\partial}{\partial y} \end{bmatrix} \begin{bmatrix} u \\ v \\ w \end{bmatrix} = \begin{bmatrix} \varepsilon_{x} \\ \varepsilon_{y} \\ \varepsilon_{z} \\ \gamma_{xy} \\ \gamma_{zy} \\ \gamma_{zx} \end{bmatrix} = [B][u]$$

$$(10)$$

应力矩阵为

$$\sigma = [\sigma_{x} \ \sigma_{y} \ \sigma_{z} \ \tau_{yy} \ \tau_{zy} \ \tau_{zx}] = D_{\varepsilon}$$
 (11)

其中, 三维本构矩阵为

设求解域上的体分布力为f,面分布力为q,集中力为 $F_i$ . 则某一瞬时单元的总能量泛函为:

$$\Pi_{P} = \int_{\Omega} \frac{1}{2} \varepsilon^{\mathsf{T}} \sigma d \, v + \int_{\Omega} c \dot{u} \, d \, v + \int_{\Omega} \rho \, \ddot{u} d \, v - \int_{\Omega} u^{\mathsf{T}} f d \, v - \int_{s_{\Omega}} u^{\mathsf{T}} q d s - \sum_{s_{\Omega}} F_{i} u_{i}$$

$$= \frac{1}{2} u^{\mathsf{T}} \int_{A}^{B} B^{\mathsf{T}} D B u d \, v + u^{\mathsf{T}} C \dot{u} + u^{\mathsf{T}} M \ddot{u} - u^{\mathsf{T}} F$$

$$= \frac{1}{2} u^{\mathsf{T}} K u + u^{\mathsf{T}} C \dot{u} + u^{\mathsf{T}} M \ddot{u} - u^{\mathsf{T}} F$$
(13)

### 1.2.2 运动微分方程和小波单元刚度

对于拱形结构来说,要对 x, y, z 方向位移分别独立插值,分别令

$$\begin{cases} u = N^e u^e = \Phi T^e u^e \\ v = N^e v^e = \Phi T^e v^e \end{cases}$$

$$w = N^e w^e = \Phi T^e w^e$$
(14)

将式 (7)、(9) 代入 (10) 则: 
$$B = \begin{bmatrix} \frac{1}{l_{ex}} \frac{\partial \Phi}{\partial \xi} \\ \frac{1}{l_{ey}} \frac{\partial \Phi}{\partial \eta} \\ \frac{1}{l_{ex}} \frac{\partial \Phi}{\partial \zeta} \end{bmatrix} T^{e}$$
 (15)

将单元求解域映射到标准求解域 $\Omega_{(0,1)}$ ,将(4)、(10)(12)(15)代入(13),由瞬时变分原理 $^{[8]}$ 令 $\delta\Pi_p=0$  $K\delta u + C\delta \dot{u} + M\delta \ddot{u} - F = 0$ 可以得到: (16)

其中:
$$K^{e} = l_{ex}l_{ey}l_{ez} \int_{\Omega_{(0,1)}} B^{T}DB \,\mathrm{d}\,\xi \,\mathrm{d}\,\eta \,\mathrm{d}\,\zeta \tag{17}$$

$$F^{e} = l_{ex}l_{ey}l_{ez} \left( \int_{\Omega_{(0,1)}} N^{e^{T}} f \, d\xi \, d\eta \, d\zeta + \int_{s_{\Omega}} qN^{e^{T}} \, d\xi \, d\eta \, d\zeta \right) + \sum_{i} N^{e^{T}} F_{i}$$

$$M^{e} = l_{ex}l_{ey}l_{ez} \int_{\Omega_{(0,1)}} \rho N^{e^{T}} N^{e} \, d\xi \, d\eta \, d\zeta$$
(18)

$$M^{e} = l_{ex}l_{ey}l_{ez} \int_{\Omega_{(0,1)}} \rho N^{eT} N^{e} \,\mathrm{d}\xi \,\mathrm{d}\eta \,\mathrm{d}\zeta \tag{19}$$

$$C^{e} = l_{ex} l_{ey} l_{ez} \int_{\Omega_{(0,1)}} cN^{eT} N \,\mathrm{d} \, \xi \,\mathrm{d} \, \eta \,\mathrm{d} \, \zeta$$
 (20)

其中:

$$K^{e} = \begin{bmatrix} K^{e,11} & K^{e,12} & K^{e,13} \\ K^{e,21} & K^{e,22} & K^{e,23} \\ K^{e,31} & K^{e,32} & K^{e,33} \end{bmatrix}$$
(21)

经计算可以得到刚度矩阵每一元素的工程显式如下

$$\begin{split} k^{e,11} &= l_{ex} l_{ey} l_{ez} \, \frac{E T^{e^{\mathrm{T}}}}{(1+\nu)(1-2\nu)} \int\limits_{\Omega_{(0,1)}} \left[ \begin{array}{c} \frac{1}{l_{ex}} \frac{\partial^2 \Phi}{\partial \xi^2} (1-\nu) + \frac{1}{l_{ey}} \frac{\partial^2 \Phi}{\partial \eta^2} \frac{1-2\nu}{2} + \frac{1}{l_{ez}} \frac{\partial^2 \Phi}{\partial \zeta^2} \frac{1-2\nu}{2} \end{array} \right] \mathrm{d}\xi \, \mathrm{d}\eta \, \mathrm{d}\zeta T^e \\ k^{e,12} &= \frac{l_{ex} l_{ey} l_{ez} E T^{e^{\mathrm{T}}}}{(1+\nu)(1-2\nu)} \int\limits_{\Omega_{(0,1)}} \left[ \begin{array}{c} \frac{1}{l_{ex}} \frac{\partial}{\partial \xi} \frac{\partial \Phi}{\partial \eta} \nu + \frac{1}{l_{ex}} \frac{1}{l_{ey}} \frac{\partial \Phi}{\partial \xi} \frac{\partial \Phi}{\partial \eta} \frac{1-2\nu}{2} \end{array} \right] \mathrm{d}\xi \, \mathrm{d}\eta \, \mathrm{d}\zeta T^e \end{split}$$

$$k^{e,13} = \frac{l_{ey}ET^{e^{\mathrm{T}}}}{(1+\nu)(1-2\nu)} \int_{\Omega_{(0,1)}} \frac{\partial \Phi}{\partial \xi} \frac{\partial \Phi}{\partial \zeta} \frac{1-2\nu}{2} d\xi \,\mathrm{d}\eta \,\mathrm{d}\zeta T^{e}$$

$$k^{e,21} = \frac{l_{ex}l_{ey}l_{ez}ET^{e^{\mathrm{T}}}}{(1+\nu)(1-2\nu)} \int_{\Omega_{(0,1)}} \left[ \frac{1}{l_{ex}} \frac{1}{l_{ey}} \frac{\partial \Phi}{\partial \xi} \frac{\partial \Phi}{\partial \eta} \nu + \frac{1}{l_{ex}} \frac{1}{l_{ey}} \frac{\partial \Phi}{\partial \xi} \frac{\partial \Phi}{\partial \eta} \frac{1-2\nu}{2} \right] \mathrm{d}\xi \, \mathrm{d}\eta \, \mathrm{d}\zeta T^{e}$$

$$k^{e,22} = \frac{l_{ex}l_{ey}l_{ez}ET^{e^{T}}}{(1+\nu)(1-2\nu)} \int_{\Omega_{(0,1)}} \left[ \frac{1}{l_{ey}^{2}} \frac{\partial^{2} \Phi}{\partial \eta^{2}} (1-\nu) + \frac{1}{l_{ey}^{2}} \frac{\partial^{2} \Phi}{\partial \xi^{2}} \frac{1-2\nu}{2} + \frac{1}{l_{ez}^{2}} \frac{\partial^{2} \Phi}{\partial \zeta^{2}} \frac{1-2\nu}{2} \right] d\xi d\eta d\zeta T^{e}$$

$$k^{e,23} = \frac{l_{ex}ET^{e^{T}}}{(1+\nu)(1-2\nu)} \int_{\Omega_{(0,1)}} \frac{\partial \Phi}{\partial \eta} \frac{\partial \Phi}{\partial \zeta} \frac{1-2\nu}{2} d\zeta d\eta d\zeta T^{e}$$

$$k^{e,31} = \frac{l_{ey}ET^{e^{T}}}{(1+\nu)(1-2\nu)} \int_{\Omega_{0,1}} \left[ \frac{\partial \Phi}{\partial \xi} \frac{\partial \Phi}{\partial \zeta} \frac{3-4\nu}{2} \right] d\xi d\eta d\zeta T^{e}$$

$$k^{e,32} = \frac{l_{ex}ET^{e^{\mathrm{T}}}}{2(1+\nu)(1-2\nu)} \int_{\Omega_{(0,1)}} \frac{\partial \boldsymbol{\Phi}}{\partial \eta} \frac{\partial \boldsymbol{\Phi}}{\partial \zeta} \,\mathrm{d}\zeta \,\mathrm{d}\eta \,\mathrm{d}\zeta T^{e}$$

$$k^{e,33} = \frac{l_{ex}l_{ey}l_{ez}ET^{e^{T}}}{(1+\nu)(1-2\nu)} \int_{\Omega_{(0,1)}} \left[ \frac{1}{l_{ez}} \frac{1}{l_{ez}} \frac{\partial^{2}\Phi}{\partial \zeta^{2}} (1-\nu) + \frac{1}{l_{ex}} \frac{1}{l_{ex}} \frac{\partial^{2}\Phi}{\partial \xi^{2}} \frac{1-2\nu}{2} + \frac{1}{l_{ey}} \frac{1}{l_{ey}} \frac{\partial^{2}\Phi}{\partial \eta^{2}} \frac{1-2\nu}{2} \right] d\xi d\eta d\zeta T^{e}$$

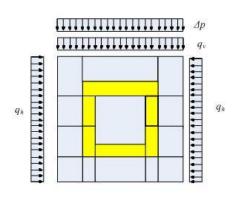


图 3 支护结构和围岩的单元划分 Fig3. The element of supporting structure and adjoining rock

式(21)同样适用于非线性分析,只要将本构矩阵换作非线性本构矩阵即可.

## 2 爆炸模拟研究

现对埋深为 30 m, 跨度为 6 m 的地下箱形硐室进行爆炸荷载作用下小波有限元模拟,直墙高度为 6 m.支护结构为 1 m 的钢筋混凝土结构,混凝土标号为 C 40,钢筋为 HRB 400 级钢筋,围岩类别为白云岩.具体参数见表 1. 炸药为 101.875 kgTNT 炸药,密度为 1 630 kg/m³,爆速为 6 717 m/s,在垂直拱顶 2m 处爆炸(见图 3). 表 1 材料参数

**Tab.1 Material parameter** 

Material	$ ho/\mathrm{kg}\cdot\mathrm{m}^{-3}$	E/GPa	ν	$F_c$ /MPa	$\mathcal{E}_u$	$F_t$ /MPa
白云岩	2 400	13.0	0.27	83	0.003	2.45
混凝土	2 400	32.5	0.27	19.1	0.002 2	2.0
钢 筋	7 850	200.0	0.27	360	6.75	360

区间B样条函数的表达式为

取硐室两边各 3 m,硐室顶部 2 m 断面,长度为 3 m 的结构进行三维模拟,用阶数 m=2,尺度 j=3 的尺度函数构造单元,由于结构对称,取一半建模,围岩和支护结构分别划分为 7 个和 4 个三维小波单元(长度方向 1 个单元,黄色部分为支护结构)。每个单元的节点数为 9 个节点,节点为等间距排列,阶数 m=2,尺度 j=3 的尺度函数有 9 个,边界小波有 2 个,内部尺度函数为 7 个,具体参见文献<sup>[1]</sup>,节点排列见图 4.

对于阻尼矩阵,采用瑞雷阻尼,岩石中的爆炸荷载取三角形荷载,由于埋深较深,基本为封闭式爆炸,假定硐室 2m 上作用均布荷载,岩石中的爆炸压力采用式(23)计算

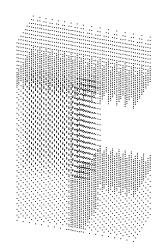


图 4 模拟模型节点排列 Fig4. The nodes arrangement of simulation model

$$P = \frac{1}{4} \rho_0 D_v^2 \tag{23}$$

其中:  $\rho_0$  为炸药密度;  $D_v$ 为炸药爆速.

此外,硐室结构还承担竖向围岩压力  $q_v$ ,水平围岩侧压力  $q_h$ ,按照普氏理论

$$q_{v} = \gamma H_{1} \tag{24}$$

其中: $\gamma$ 为岩石容重; $H_1$ 为围岩压力拱高度,岩石抗压强度为 $R_c$ ,围岩稳定性安全系数为 $f_{\leftarrow}R_c/10$ ,则压力拱高度为

$$H_1 = a_1 / f_i \tag{25}$$

$$a_1 = a + ftg (45^{\circ} - \frac{\varphi}{2})$$
 (26)

式中: a为毛洞半跨; f为拱高. 由于埋深较大,洞高较小,侧向压力按均布荷载计算

$$q_h = qtg^2 (45^\circ - \frac{\varphi}{2}) \tag{27}$$

混凝土采用整体式模型,钢筋配筋率为 $\rho_x=0.10$ , $\rho_y=0.10$ , $\rho_z=0.06$ ,钢筋混凝土本构矩阵为:

$$D = D_c + D_s \tag{28}$$

其中 $D_c$ 为混凝土的本构矩阵,按式(12)计算. $D_c$ 为钢筋的本构矩阵,按式(29)计算.

运动方程的解法采用增量法<sup>[9]</sup>,经过 Matlab 编程<sup>[10]</sup>计算得到的顶板位移和 Y 向应力,与 ANSYS/LS-DINA 模拟<sup>[11]</sup>得到结果对比如图 5 和表 2 所示.

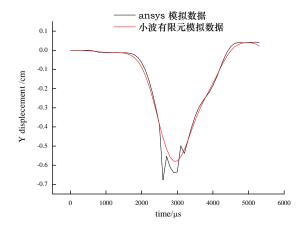


图 5 顶板中部竖向位移 Fig.5 The vertical displacement middle of roof

## 表 2 支护顶板节点竖向最大应力对比(MPa)

Tab.2 The vertical stress of roof (MPa)

模拟方法	节点号									
	1 支座	2	3	4	5	6	7	8	9 中点	
小波模拟	9.3	10.7	12.4	13.5	14.9	15.7	16.6	17.2	17.8	
Ansys 模拟	9.2	10.6	12.5	13.5	14.8	15.6	16.5	17.1	17.5	

表 2 中数据为板底由支座到中点 9 个节点处竖向最大拉应力. 由图 5 和表 2 可见,区间样条小波模拟结果与 7 168 个传统有限元的数值模拟结果基本一致,尤其是顶板中部竖向位移非常接近,并且避免了数值震荡.

## 3 结论

本文推导了小波单元的三维转换矩阵,构造了三维区间 B 样条小波单元,并用其对爆炸荷载作用地下箱形结构进行数值模拟,通过与 ANSYS/LS-DYNA 数值模拟结果对比可见:

- (1) 12 个三维小波单元与 7168 个传统有限元单元的模拟精度相当接近,并且避免了数值震荡.
- (2) 区间 B 样条小波单元的刚度矩阵是非奇异性矩阵,收敛速度快,计算效率提高 50%以上.

以上研究表明构造的三维 B 区间样条小波单元是正确的,小波有限元可以应用于爆炸等动载作用下结构动力响应的数值模拟.

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(下转第832页)

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## Numerical analysis of collision between truck and coal transporting trestle

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**Abstract:** In order to improve the capacity of collision of trestle structure, the mechanism of collision between truck and trestle structure was studied. Based on actual accident investigations, a classical collision accident of a coal transporting trestle hit by vehicle was simulated. The failure modes, stress and displacement response of the castellated beam and the bracket were analyzed. Results indicate that under impact load the stress of castellated beam was maximal at the supporting regions and impact site. Bracket has the feature of tension-torsion. The computational results can be used for further studies on the truck- trestle impact problem.

Key words: coal transporting trestle; castellated beam; bracket; collision

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(上接第821页)

# The three-dimension spline wavelet finite element Construction and application on antiknock of underground box structure

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Abstract: Because of the characteristics of special material and high-speed loading becomes it easy for numerical oscillation to occur in traditional finite element simulation when the underground structure is subject to blasting load. The division of elements is numerous, and the computation inefficient. The wavelets finite element has the characteristics of many dimensions and many resolutions. And the precision can be improved by increasing yardstick of wavelets functions or measurement functions. The three-dimension transition matrix was derived. The new element of three-dimension interval B-spline was constructed using the measurement functions as interpolation function, which was applied on antiknock of underground box structure under blasting load through programming with Matlab software. The result shows that the computation precision is high with little elements and the computation efficiency is improved comprised with common finite element.

Key words: blast load; underground box structure; dynamic action; interval B-spline wavelet; element construction

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